

Molecular Hydrogen for RAMSES

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Outline

1. An introduction to H_2
2. Ramses-RT with H_2
3. Full galaxy simulation

H₂ in Galaxies

The Circinus molecular cloud with new star and protostellar jet

- Most common molecule in the Universe
- Important to cooling the interstellar medium
- No dipole moment and is light, leading to high emission temperature
- Difficult to detect directly, usually inferred from CO by conversion factor
- Dense, cold molecular clouds are linked to star formation

Molecular Cloud Schematics

Lyman
Werner
radiation
(11.20-13.60 eV)

Ionizing
radiation
(≥ 13.60 eV for H I,
 ≥ 15.42 eV for H₂)

easily observable?



H II gas



H I gas



H₂ gas



CO gas



HCN gas

hot

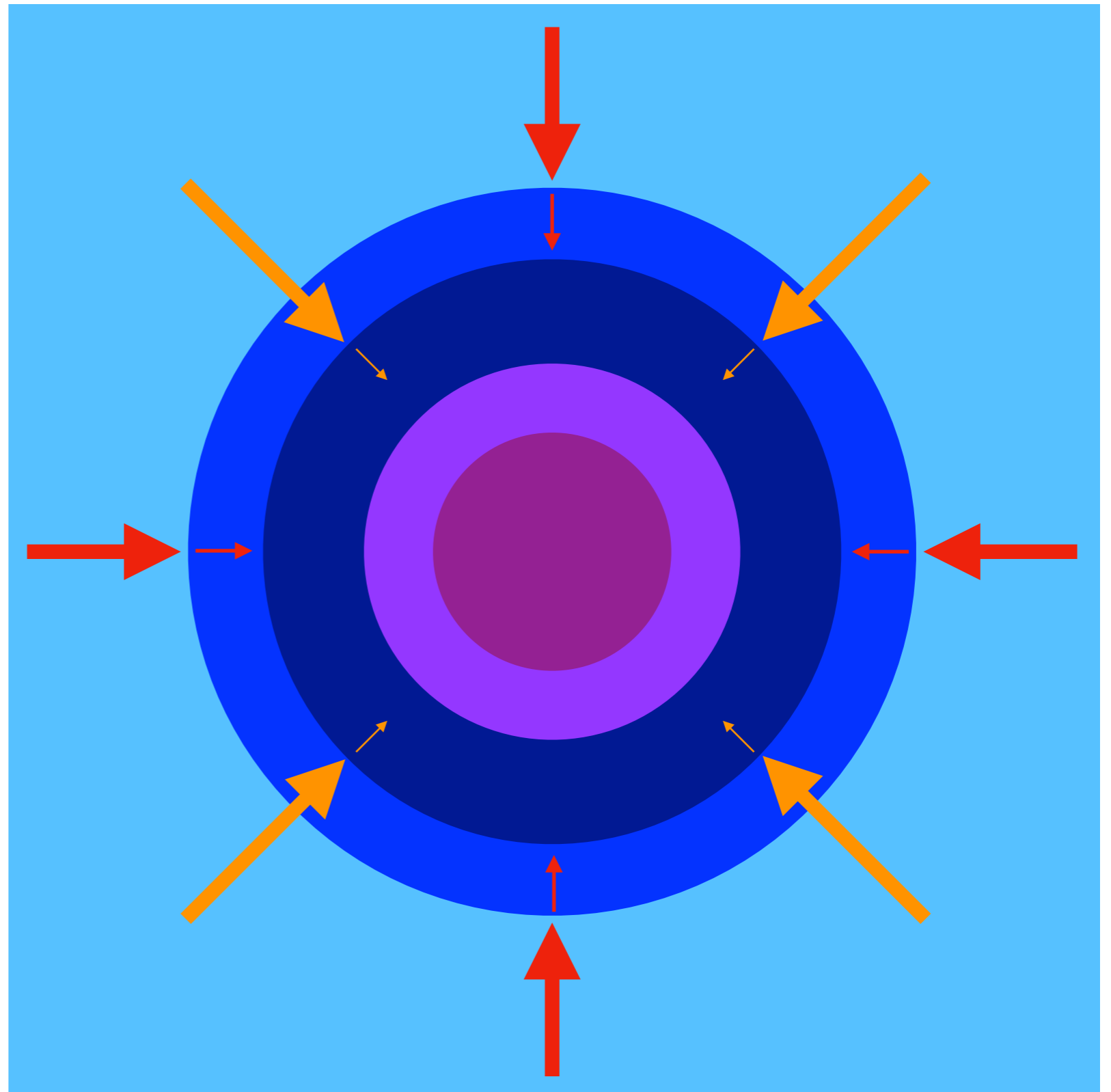
diffuse



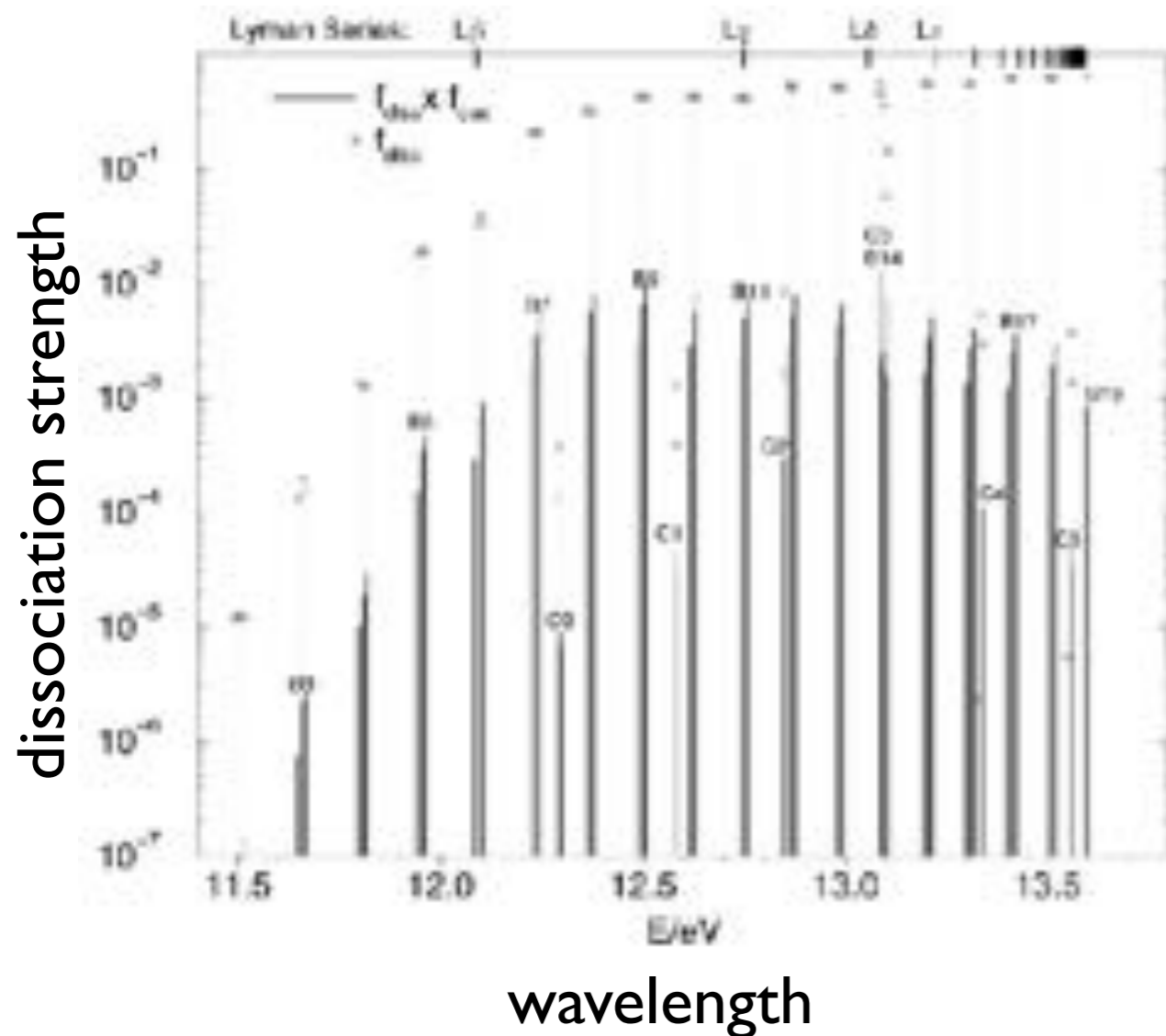
cold



dense



H₂ Self-Shielding



Source: Haiman et al. 2000

- Not every absorption of a Lyman-Werner photon (11.2 to 13.6 eV) leads to dissociation
- Different lines of Lyman-Werner photons dissociate H₂ at different rates
- Stronger lines absorbed at cloud surface, weaker lines sneak deeper
- Heisenberg uncertainty means that finite molecular excitation time leads to natural line width
- The strong lines interfere with the weak lines
- Deeper into H₂ cloud, absorption becomes harder

Ramses-RT

- Rosdahl et al. 2013
- Moment-based radiative transfer scheme for RAMSES
- Computation independent of source number
- Non-equilibrium chemistry of HI, HII, HeI, HeII, and HeIII coupled to ionising photon groups and thermal state of gas



Adding H₂ to RAMSES-RT

$$\begin{aligned} \dot{n}_{\text{H}_2} = & \underbrace{\alpha_{\text{H}_2}^{\text{Z}}(T) Z f_{\text{d}} n_{\text{H}} n_{\text{H}_1}}_{\text{Formation on dust}} + \underbrace{\alpha_{\text{H}_2}^{\text{GP}}(T) n_{\text{H}_1} n_{\text{e}}}_{\text{Gas-phase formation}} \\ & + \underbrace{\beta_{3\text{B}}(T) n_{\text{H}_1}^2 (n_{\text{H}_1} + n_{\text{H}_2}/8)}_{\text{Formation by three-body collisions}} \\ & - \underbrace{\beta_{\text{H}_2\text{H}_1}(T) n_{\text{H}_1} n_{\text{H}_2}}_{\text{Collisional destruction with HI}} - \underbrace{\beta_{\text{H}_2\text{H}_2}(T) n_{\text{H}_2}^2}_{\text{Collisional destruction with itself}} \\ & - \underbrace{\Gamma_{\text{H}_2}^{\text{LW}}(N_{\text{H}_2}) n_{\text{H}_2}}_{\text{Photodissociation by Lyman-Werner photons}} - \underbrace{\Gamma_{\text{H}_2}^+(N_{\text{H}_1}) n_{\text{H}_2}}_{\text{Photoionisation by ionising photons}} - \underbrace{\xi_{\text{H}_2} n_{\text{H}_2}}_{\text{Cosmic ray ionisation}} \end{aligned}$$

Formation on dust (Hollenbach & McKee 1979; Jura 1974; Gry et al. 2002; Habart et al. 2004)

Gas-phase formation (McKee & Krumholz (2012)

Formation by three-body collisions (Forrey 2013; Palla et al. 1983)

Collisional destruction with HI (Dove & Mandy 1986)

Collisional destruction with itself (Martin et al. 1998)

Photodissociation by Lyman-Werner photons (Sternberg et al. 2014)

Photoionisation by ionising photons (Abel et al. 1997)

Cosmic ray ionisation (Indriolo & McCall 2012; Gong et al. 2017; Glassgold & Langer 1974)

Additional Thermal Processes

$$\mathcal{H} = \sum_j^{\text{H}_2, \text{H I}, \text{He I}, \text{He II}} n_j \sum_{i=1}^M c_i N_i (\bar{\epsilon}_i \sigma_{ij}^E - \epsilon_j \sigma_{ij}^N) + \mathcal{H}_{\text{PE}}(T) + \mathcal{H}_{\text{UVP}}(T) + \mathcal{H}_{\text{H}_2}(T) + \mathcal{H}_{\text{CR}}(T)$$

(heating)

$$\mathcal{L} = [\zeta_{\text{H I}}(T) + \psi_{\text{H I}}(T)] n_e n_{\text{H I}} + \zeta_{\text{He I}}(T) n_e n_{\text{He I}} + [\zeta_{\text{He II}}(T) + \psi_{\text{He II}}(T) + \eta_{\text{He II}}^{\Lambda}(T) + \omega_{\text{He II}}(T)] n_e n_{\text{He II}} + \eta_{\text{H I}}^{\Lambda}(T) n_e n_{\text{H I}} + \eta_{\text{He II}}^{\Lambda}(T) n_e n_{\text{He II}} + \theta(T) n_e (n_{\text{H I}} + n_{\text{He I}} + 4n_{\text{He II}}) + \varpi(T) n_e + \Lambda_Z(T) + \Lambda_{\text{H}_2}(T)$$

(cooling)

Photoelectric effect heating (Bakes & Tielens 1994, Wolfire et al. 2003)

Heat from UV pumping (Baczynski 2015, Draine & Bertoldi 1996, Burton et al. 1990)

H₂ formation heating (Hollenbach & McKee, Omukair 2000)

Cosmic ray ionization heating (Glassgold et al. 2012)

H₂ collisional destruction cooling (Halle & Combes 2013, Hollenbach & McKee 1979)

Previous Self-Shielding Models

- Galaxy simulations use the Draine and Bertoldi 1996 self-shielding as a function of column density
- Decreases H₂ destruction
- Problems: simulations use volume density and ray tracing is expensive for galaxies with multiple sources
- Many approximations to convert column density to volume density, local and non-local

$$f_{\text{shield}}(N_2) = \frac{0.965}{(1 + x/b_5)^2} + \frac{0.035}{(1 + x)^{0.5}} \times \exp[-8.5 \times 10^{-4}(1 + x)^{0.5}]$$

$x = N_2 / 10^{14}$
 b is a constant

Our Self-Shielding Model

- Takes advantage of moment-based radiative transfer that treats photons like a fluid
- Instead of decreasing H₂ destruction, we enhance Lyman-Werner destruction
- Use a constant factor, S ; calibrated with Bialy et al. 2016 HI-H₂ transition depth
- Cumulative destruction of LW photons while they travel through each cell is similar

$$D_{LW} = c\sigma_{LW}n_{H_2} \rightarrow S c\sigma_{LW}n_{H_2}$$

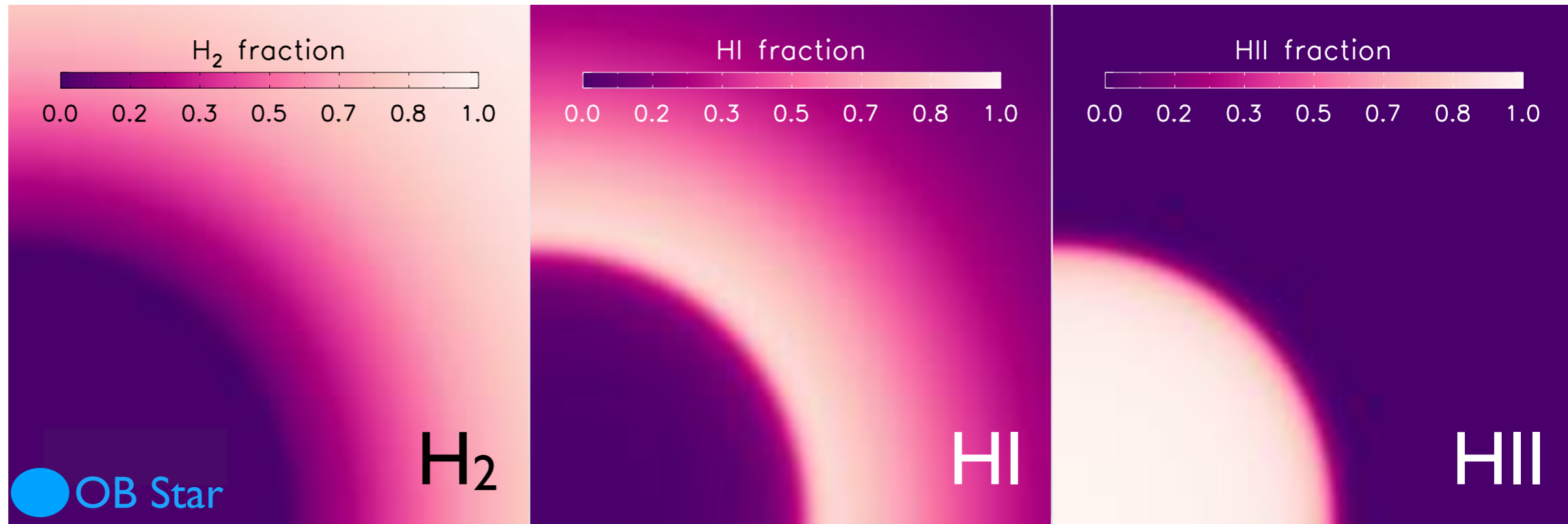
Molecular Strömgren Sphere

Originally describes radiation from an OB star ionising an atomic medium

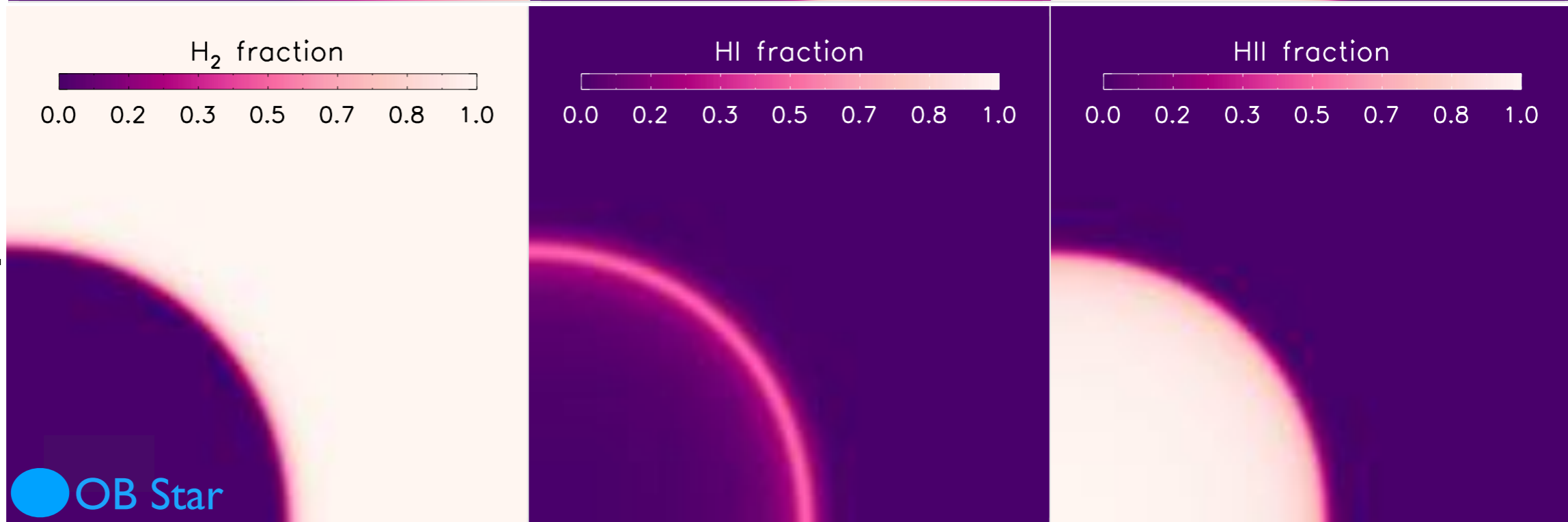
We describe ionizing and dissociating radiation in an atomic and molecular medium

Molecular and atomic fronts grow according to analytical expressions

No self-shielding



With self-shielding



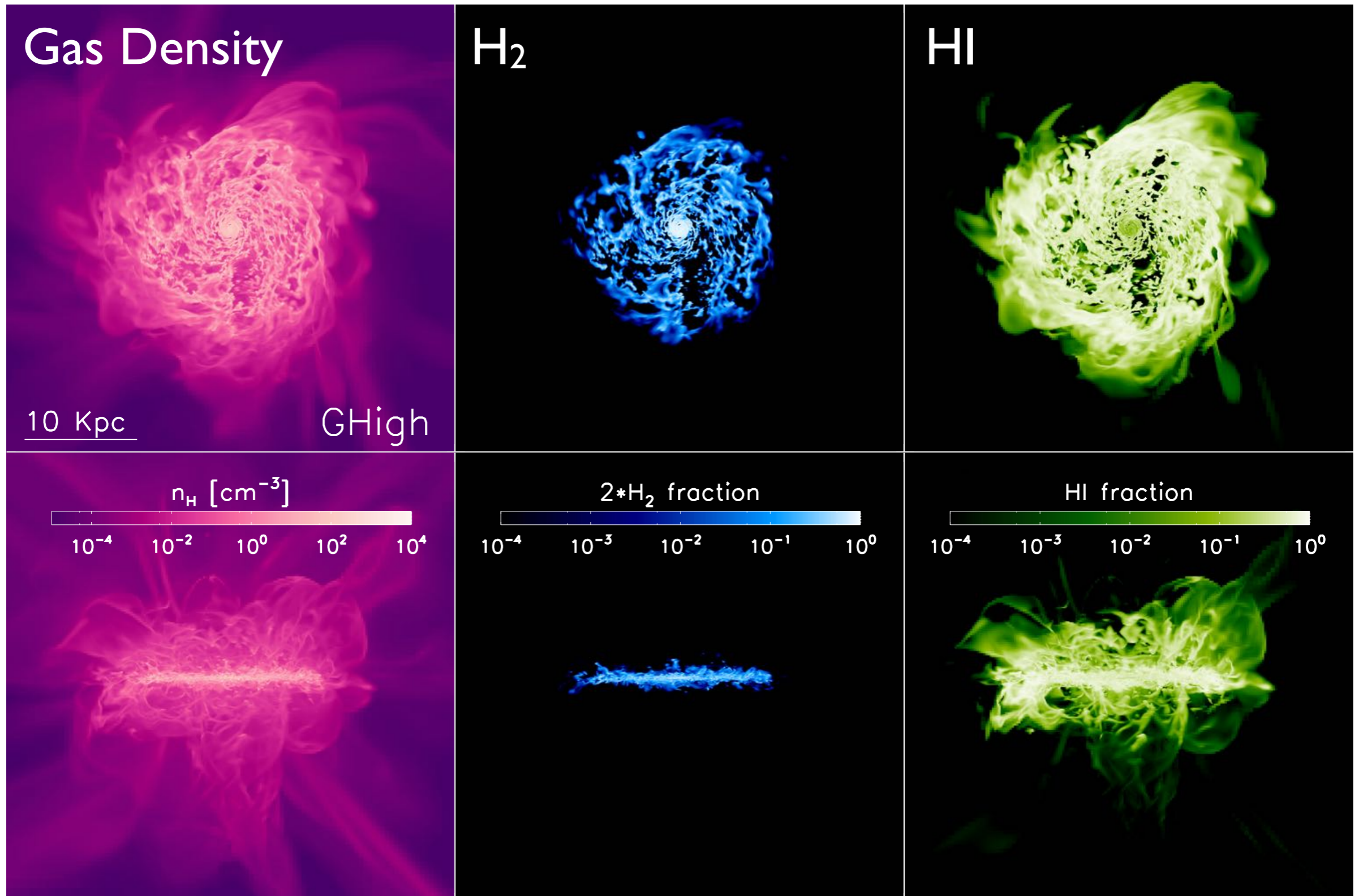
Galaxy!

Milky Way-like, isolated disc galaxy (from the AGORA project, Kim et al. 2016)

Resolution 6 pc, 400 kpc box

Star formation recipe independent of H_2 , is total gas density-based (Rasera & Teyssier 2006)

Delayed cooling thermal feedback for supernovae (Teyssier et al. 2013)



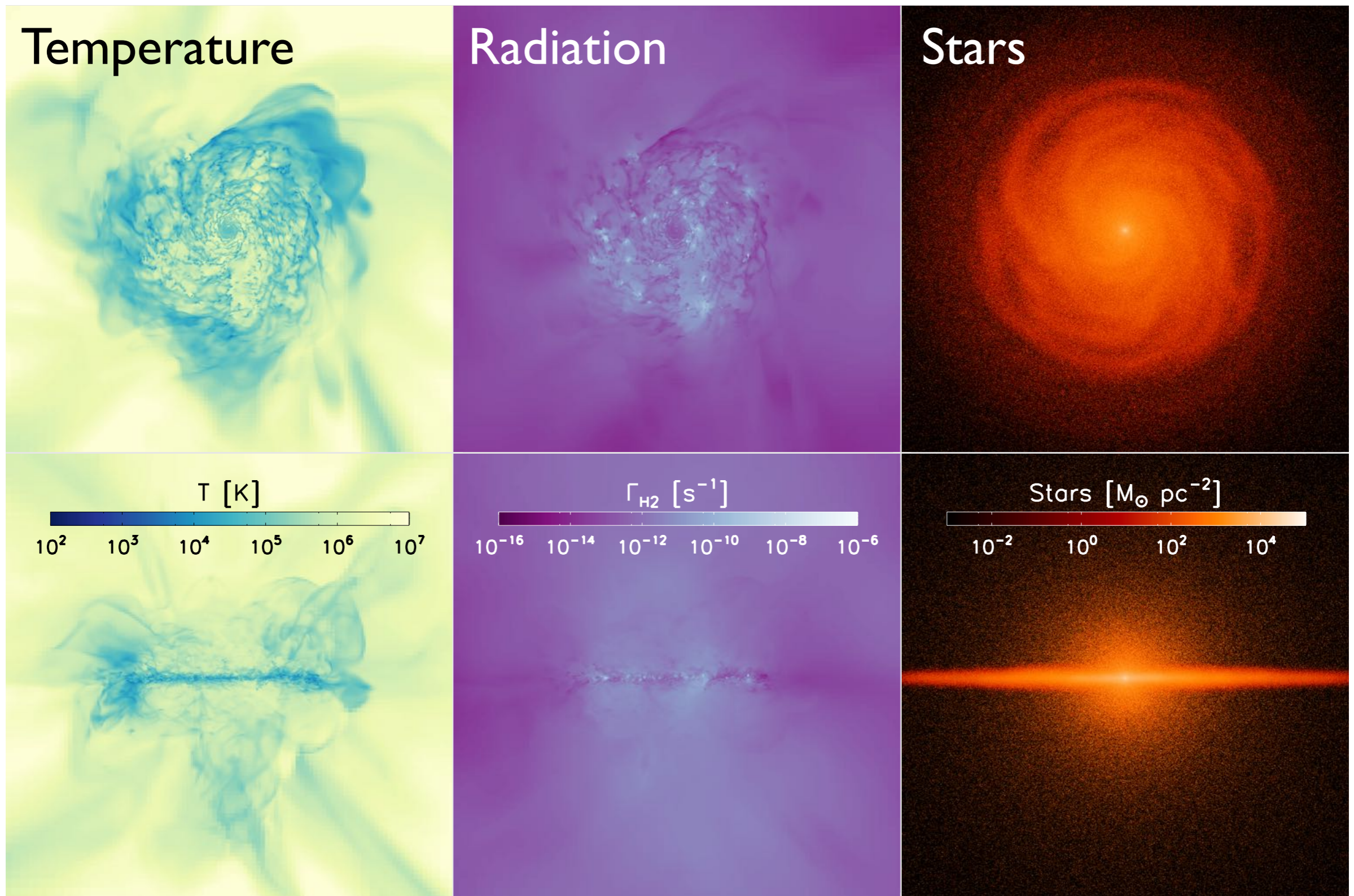
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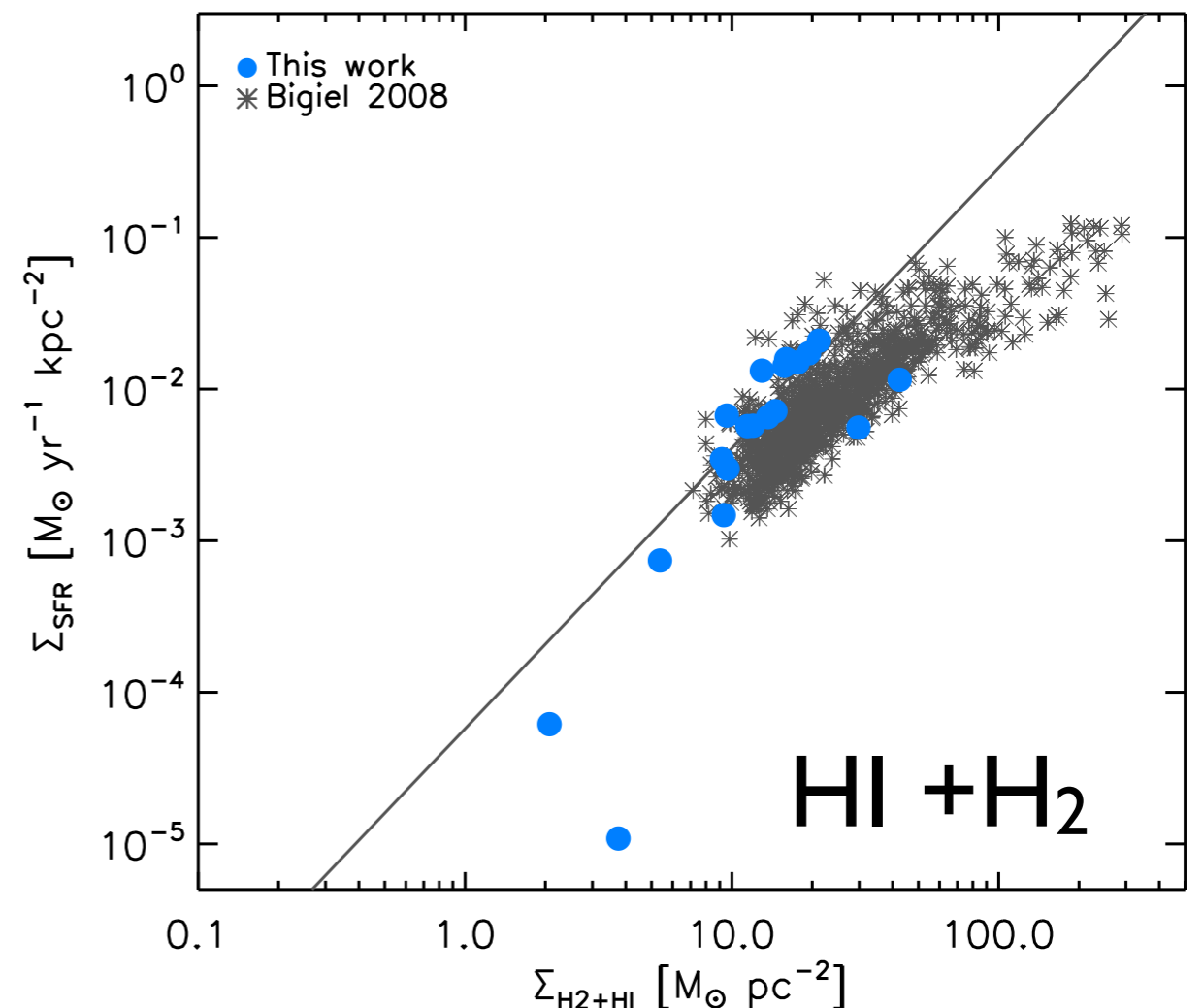
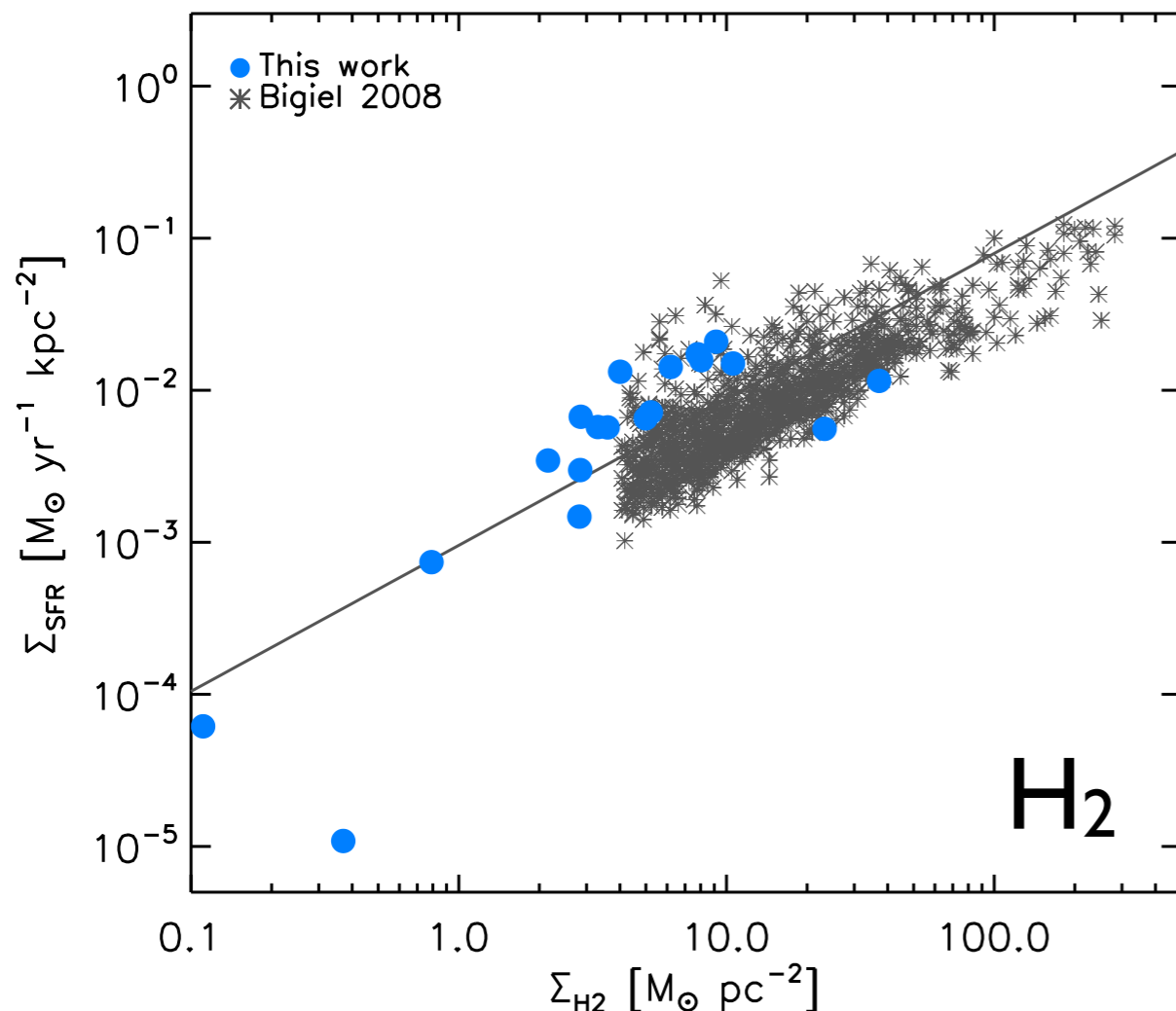
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Kenincutt-Schmidt

- Established, observed relation between neutral gas and SFR (Schmidt (1959) and Kennicutt (1998))
- More recently discovered relation between H_2 and SFR (e.g. Bigiel et al. 2008)
- We match both, aside from quenched central points
- Scruba 2011 extend this relation to lower densities, and we match fall in HI- H_2 relation



Phase Diagrams

All plots weighted by gas mass fraction

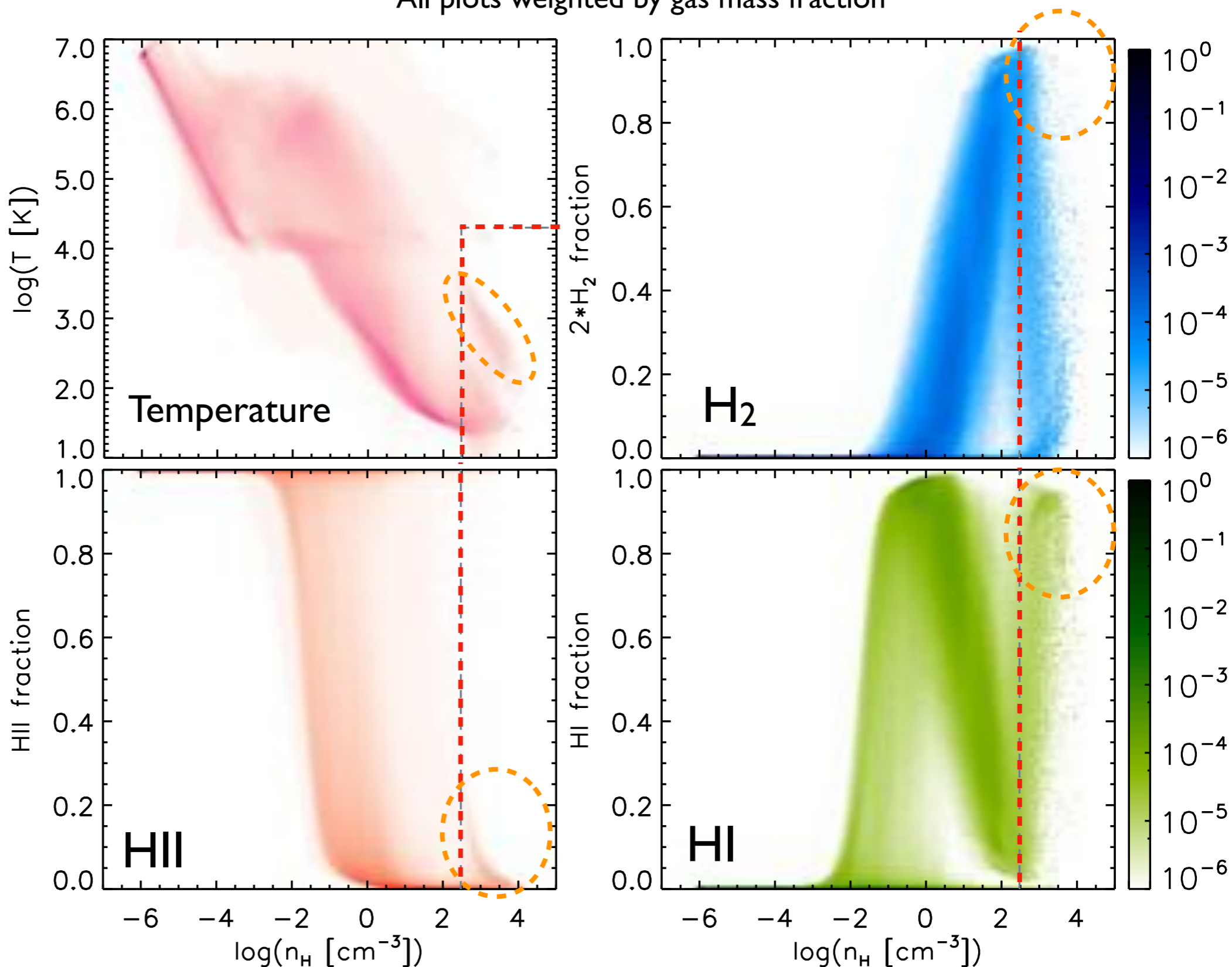
 Unresolved HII regions
 Star formation threshold

HII sparse, hot gas

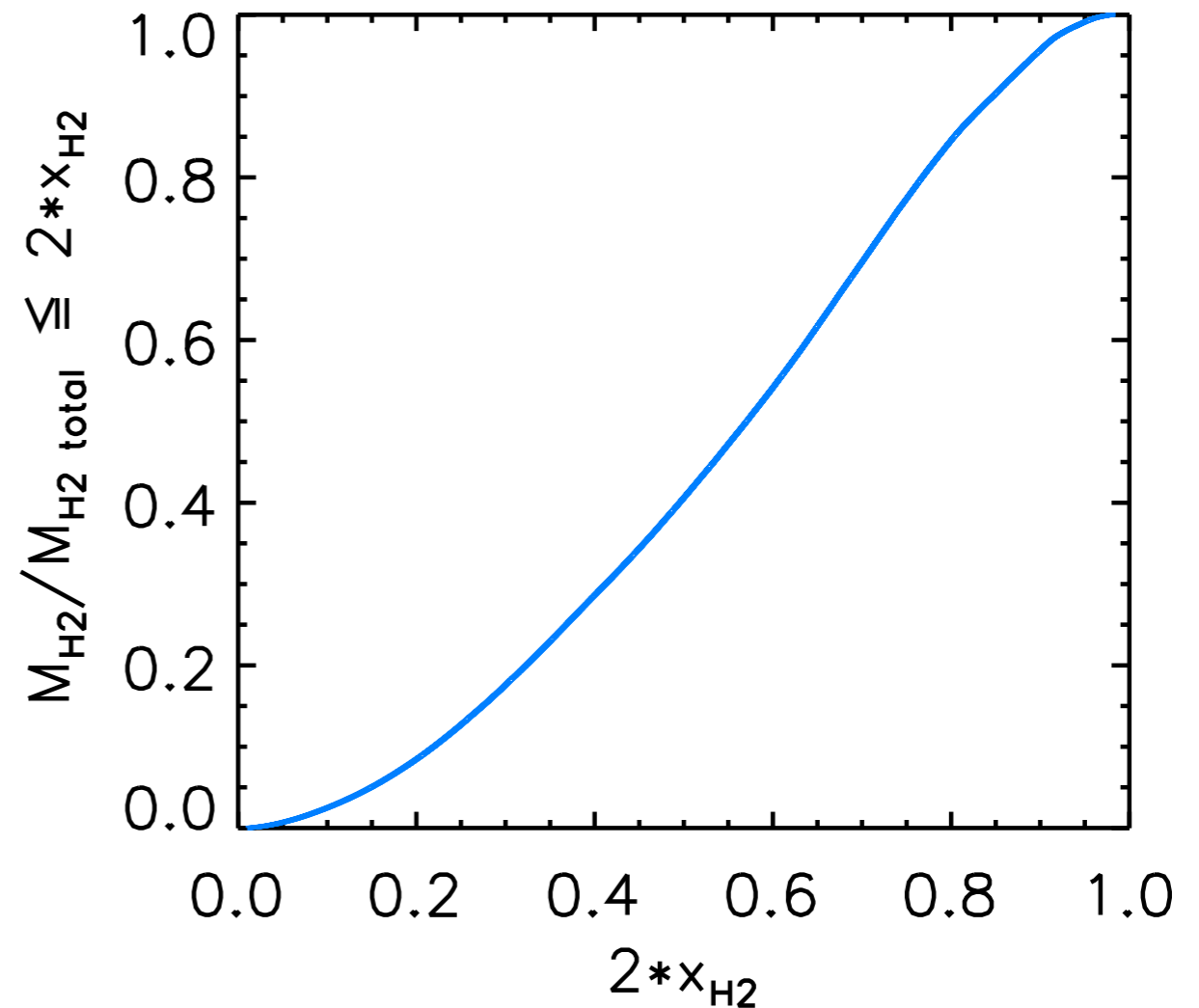
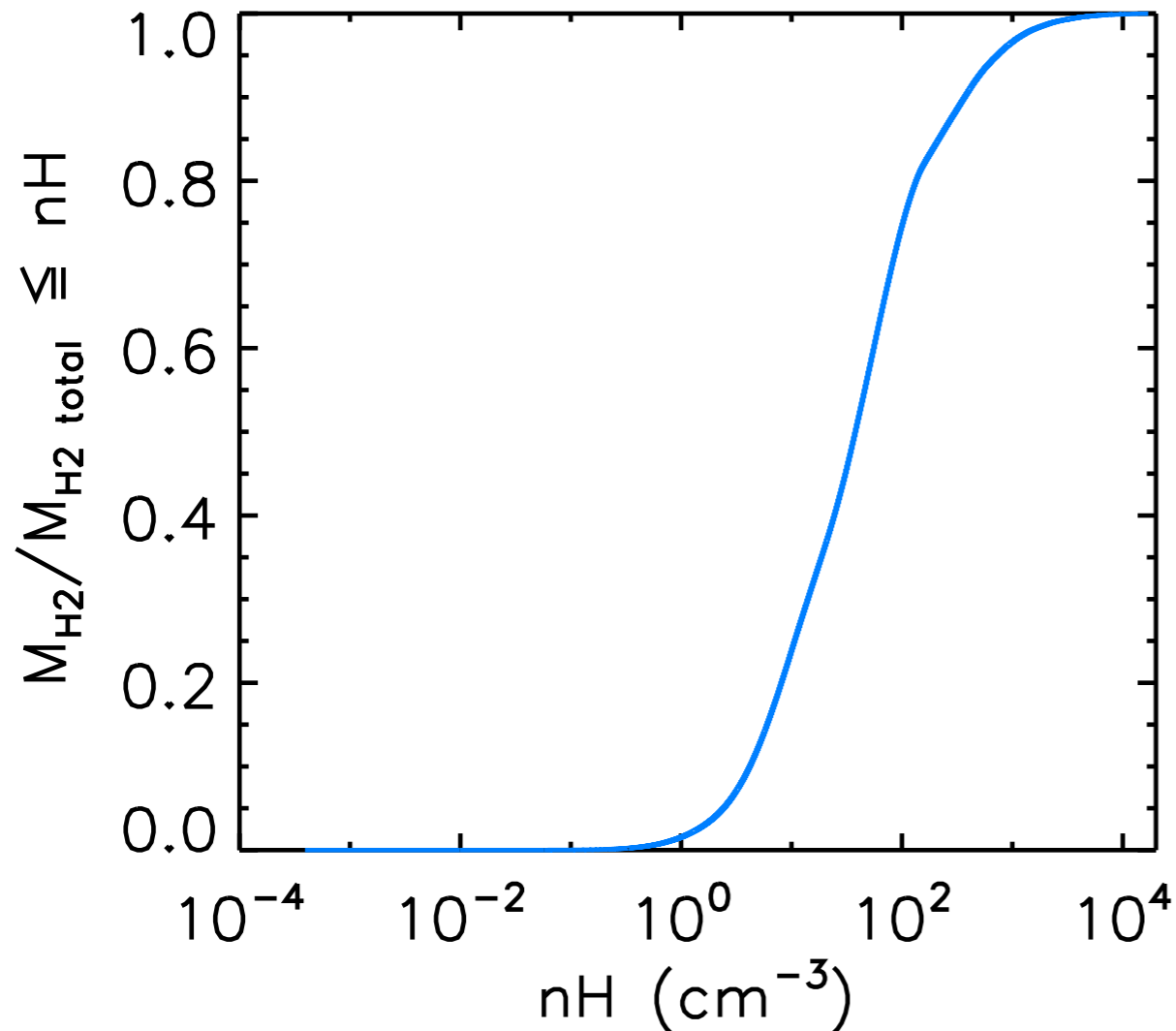
H₂ cold, dense gas

HI in between

Cannot resolve 100%
molecular regions or
HII regions around
young stars

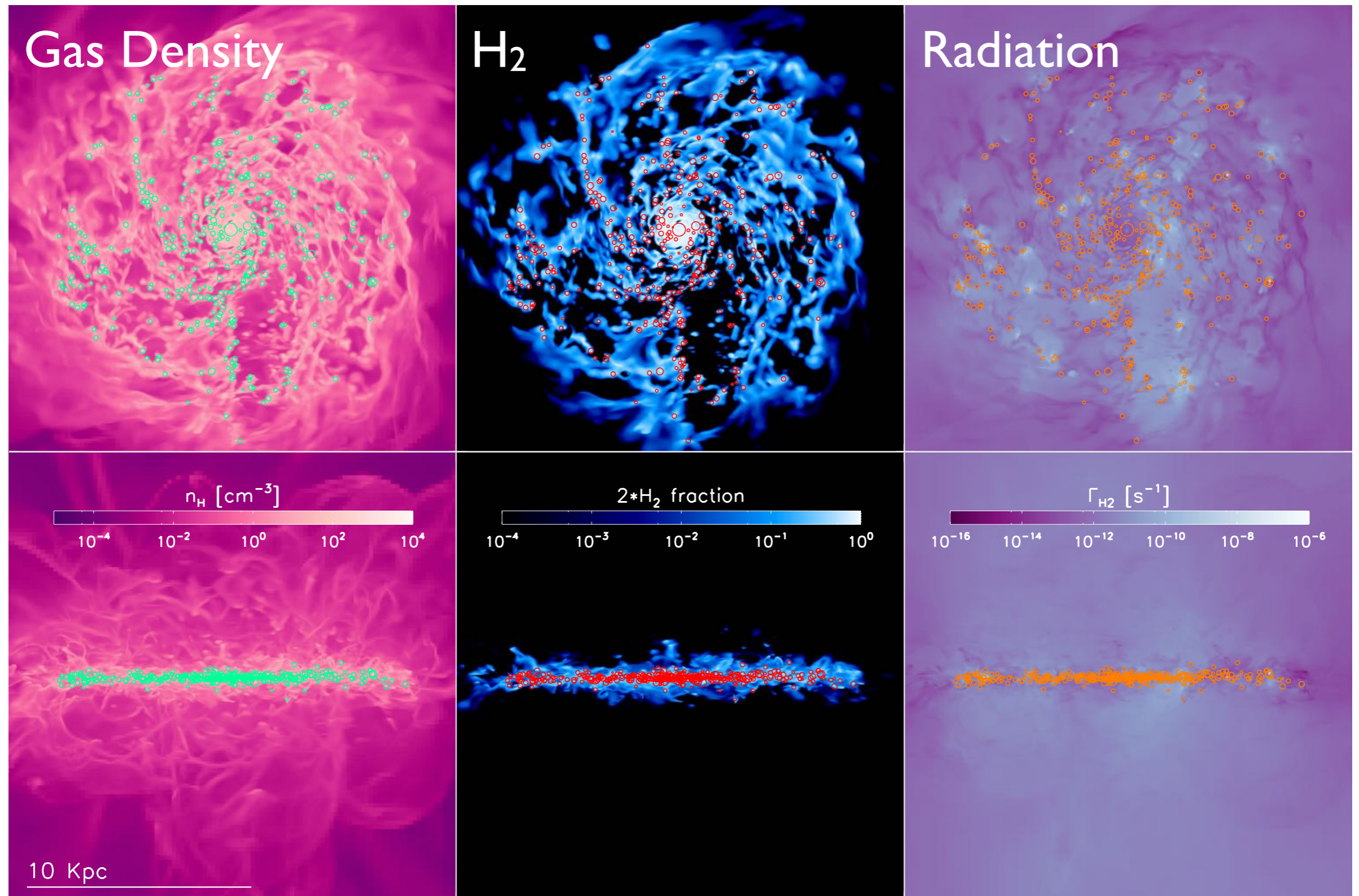


Where is the H₂?



- 50% of H₂ is at densities 37.8 cm⁻³ and lower
- 50% of H₂ is at fractions 0.57 and lower

Clump Finding With PHEW



Conclusions

- Successfully modelled non-equilibrium H_2 chemistry using local methods with Ramses-RT
- Introduced new radiative transfer-based self-shielding model for galaxy simulations
- Tested model in idealised situations against benchmark models to demonstrate robustness
- Realistic galaxy
- Much of the H_2 gas is diffuse and mixed with HI
- In progress: CO chemistry to investigate the CO- H_2 conversion factor and the existence of a “CO-dark” H_2 component of the ISM

Chemistry: Nickerson Sarah, Teyssier Romain, Rosdahl Joakim, 2018, MNRAS, 479, 3206

Galaxy: Nickerson Sarah, Teyssier Romain, Rosdahl Joakim, 2018, arXiv:1809.01657,
submitted to MNRAS

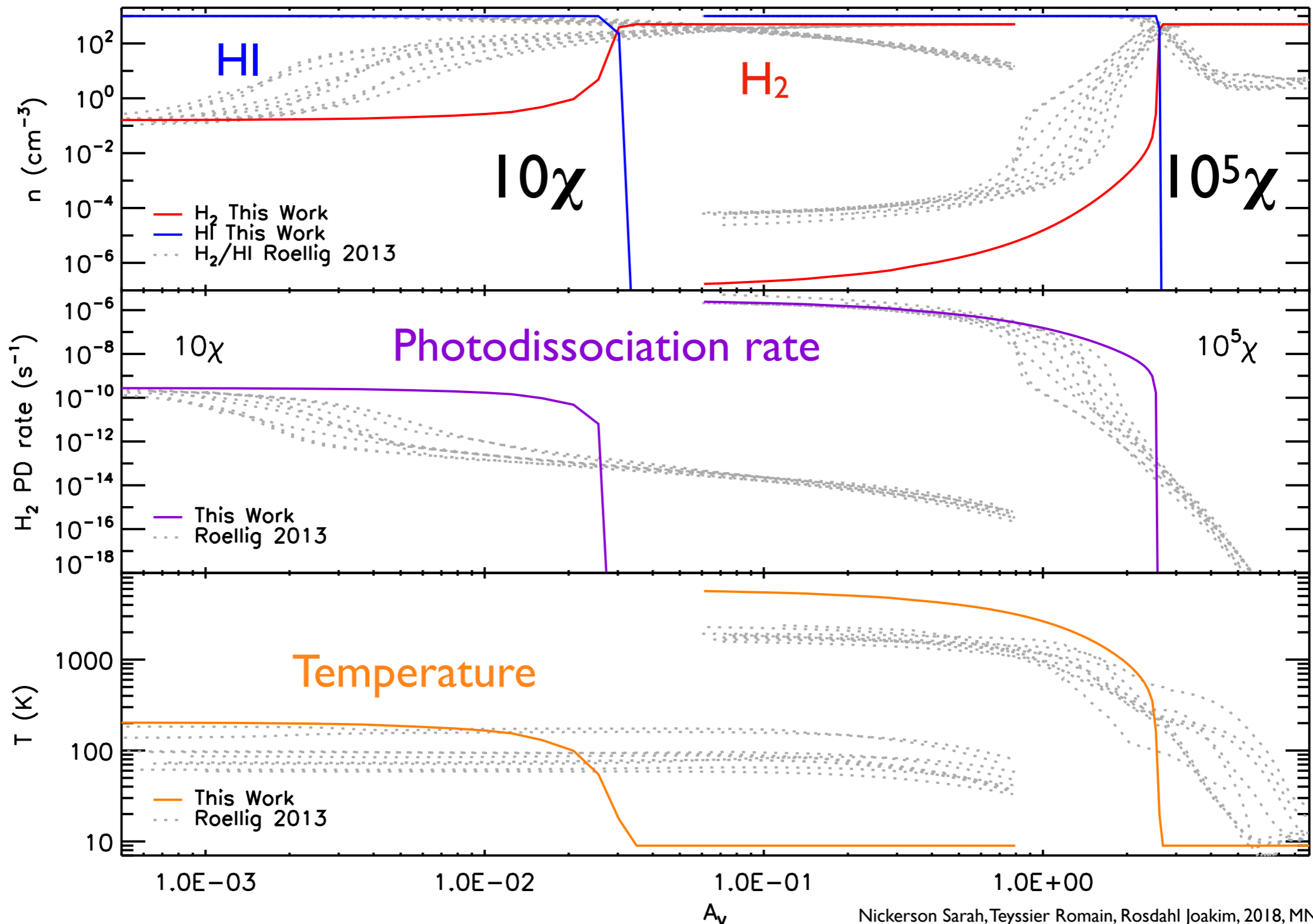
PDR Code Comparison

Photodissociation regions (PDRs) UV dominated (6 to 13.2 eV) and host most galactic atomic and molecular gas

Roellig 2013 benchmark comparison of 10 PDR codes

PDR codes are 1D so can treat self-shielding exactly

We reproduce transition point, but not shape



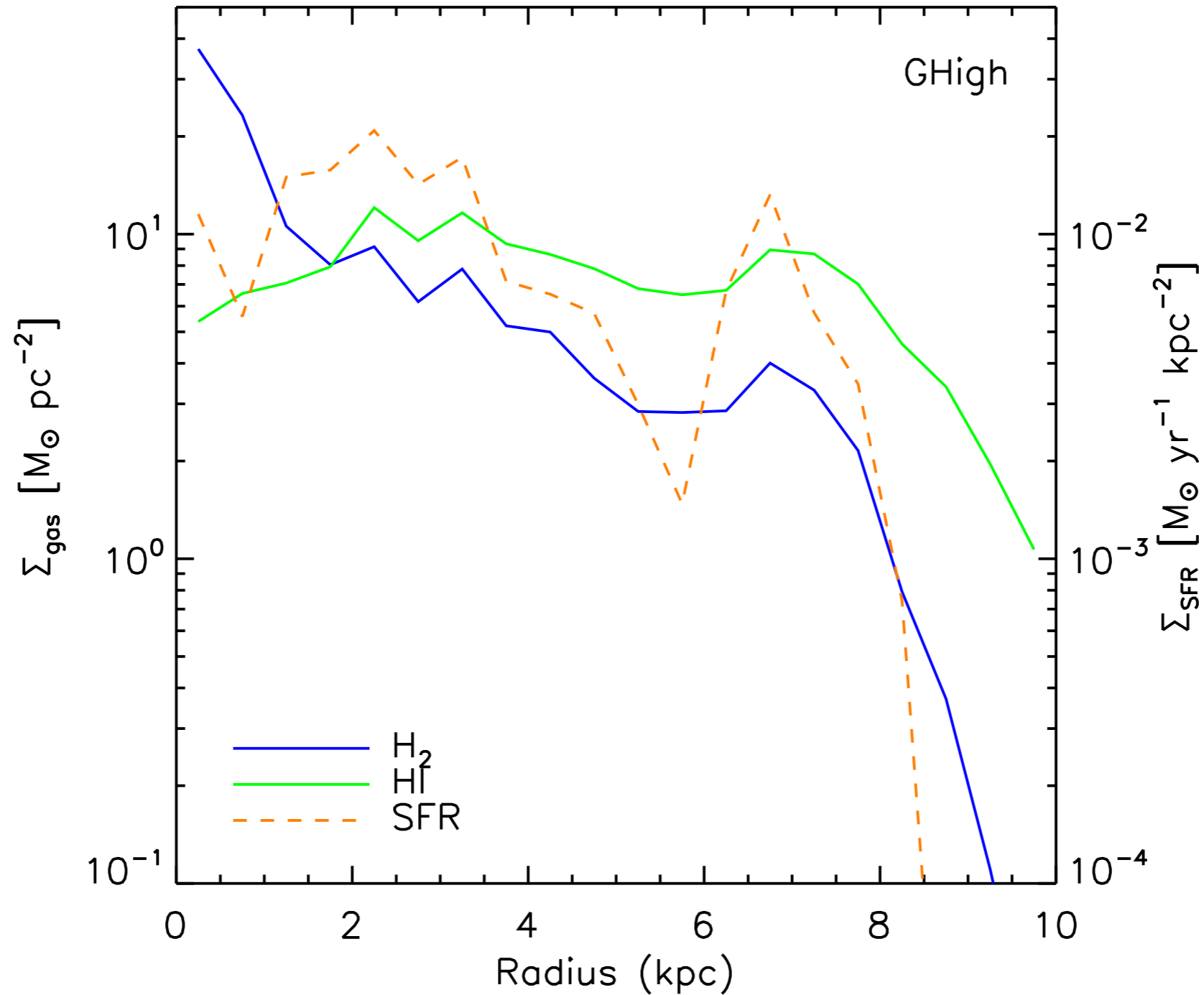
Radial Profile

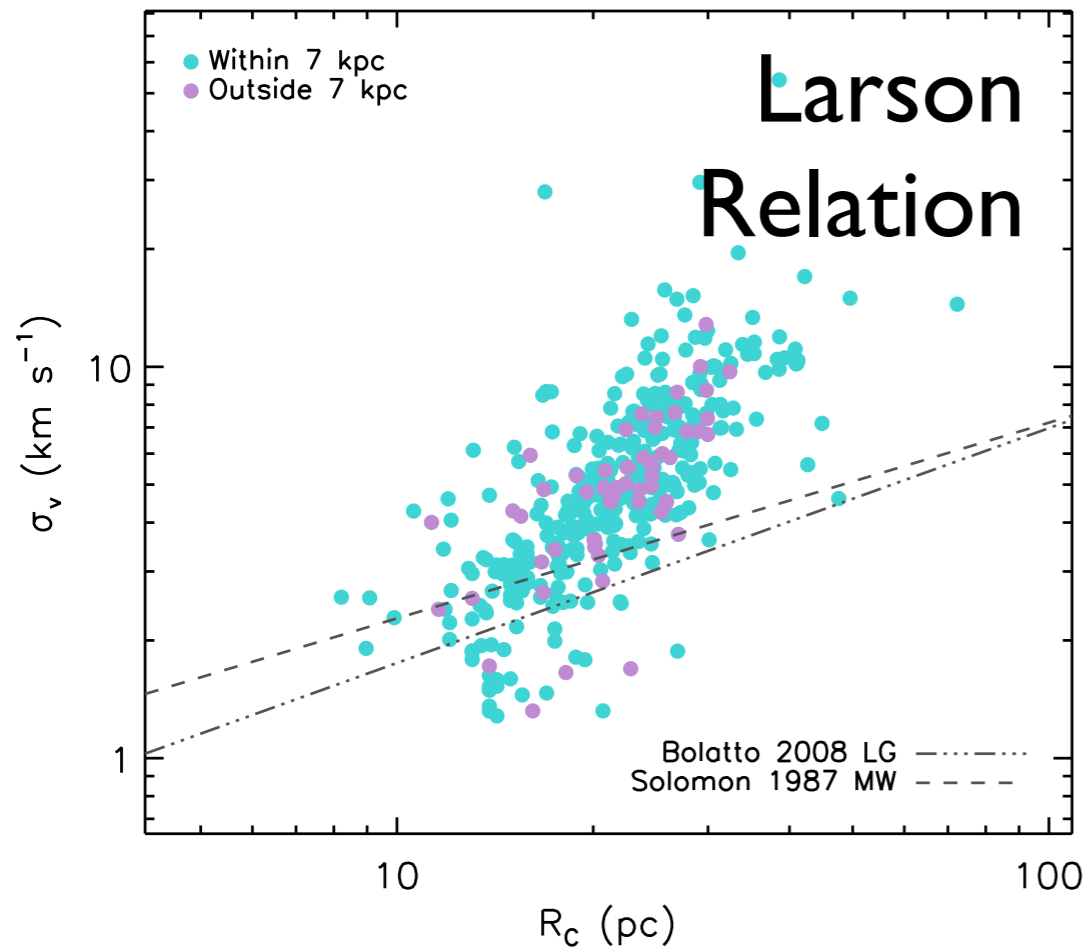
Reproduce central H₂ spike and flat HI profile

SFR trends mimic H₂ trends, except for galactic centre

Central bulge may cause morphological quenching

Surface density and SFR numbers are similar to observed galaxies around Milky Way mass





- Clouds are 96% molecular
- Slightly more at medium masses than Milky Way, fewer low mass clouds
- Inner and outer galaxy cloud populations show no difference
- Do not follow traditional Larson relation (1981), but do with newer Heyer relation (2009)

